

DN Operator

$$\begin{cases} -\Delta u = 0 \\ u|_{\partial\Omega} = f \in C^\infty? \end{cases} \quad T_{DN} f = \partial_n f$$

upper half plane

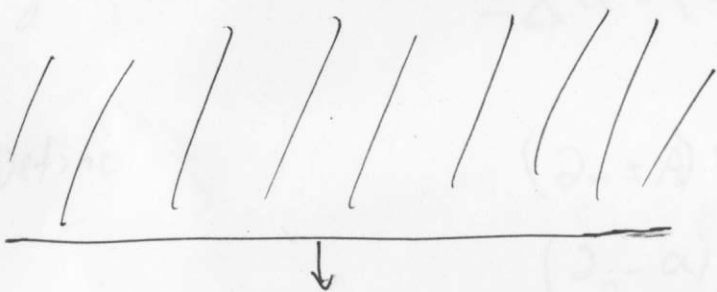
$$\begin{cases} u(x,0) = f(x), x \in \mathbb{R}^d \\ -\Delta u = 0 \end{cases}$$

claim. $T_{DN} f = (-\Delta)^{1/2} f$

• Define $Tf(x) = -u_y(x,0)$

then $-u_y(x,y)$ solves

$$\begin{cases} v(x,0) = Tf \\ -\Delta v = 0 \end{cases}$$



• $T^2 f = T(Tf) = T(-u_y(x,0)) = u_{yy}(x,0) = -\Delta_x f$

• T is positive: By Green's formula:

$$\langle Tf, f \rangle = \int_{\mathbb{R}^d} (Tf)f = \int_{\mathbb{R}^d} -u_y(x,0) f(x) = \int_{\mathbb{R}^{d+1}} |\nabla u|^2 + u \Delta u \geq 0$$

Basically true in general, symbol for $-\Delta$ is $|\xi|^2$.

On Riemannian manifold (M, g) have basically principal symbol of T_{DN} is (thesis by Jan Möllers)

$$a(x, \xi) = \sqrt{\sum_{j,k=1}^{n-1} g^{jk}(x) \xi_j \xi_k}$$

$$(Taf(x) = \int a(x, \xi) \hat{f}(\xi) e^{ix\xi} d\xi)$$

This can be seen by "factoring" the Laplacian. (or its symbol).

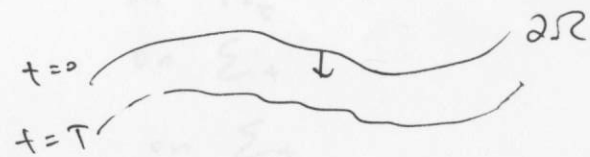
Say can factor

$$-\Delta u = (\underbrace{\partial_n - a}_{\text{tangential}}) (\underbrace{\partial_n + A}_{\text{tangential}}) u + \underbrace{Ru}_{\text{smoothing}}$$

Define

$$(\partial_n + A) u = v$$

$$(\partial_n - a) v = -Ru \in C^\infty$$



Elliptic regularity $\Rightarrow u \in C^\infty$ in \mathbb{R}^n .

$$v(T) = (\partial_n + A) u(T) \in C^\infty$$

$$\frac{dv}{dt} - av \in C^\infty$$

Solve backwards, reg. for heat eqn $\Rightarrow v(0) \in C^\infty$

Finally,

$$\partial_n f = \frac{dv}{dt}(0)$$

$$= -A u(0) - v(0)$$

$$= -A(0)u - v(0)$$

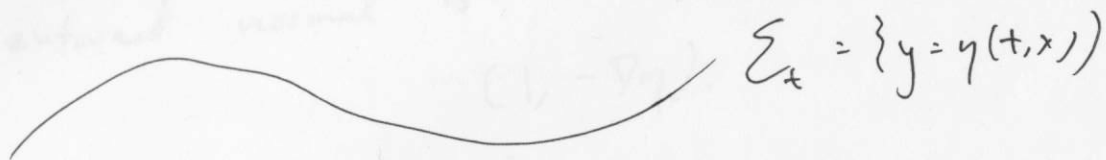
DAW for Water Wave Equation

(On WWE with Surface Tension by Alazard, Burg, Zuilly)

(3)

Setup

$$\{(t, x, y) \in [0, T) \times \mathbb{R}^d \times \mathbb{R} : (x, y) \in \Omega_t\}$$



$$\begin{cases} \Delta_x \Psi + \partial_y^2 \Psi = 0 \\ \partial_t \eta = \partial_t \Psi - \nabla \eta \cdot \nabla \Psi \\ \partial_t \Psi = -g\eta + \kappa H(\eta) - \frac{1}{2} |\nabla \Psi|^2 - \frac{1}{2} |\partial_t \Psi|^2 \\ \partial_n \Psi = 0 \end{cases}$$

in Ω_t
on Σ_t
on Σ_t
on Γ

Defn.

Let $\Psi(t, x) = \Psi(t, x, \eta(t, x)) = \Psi$

DN-operator:

$$(\partial_n \Psi) \Psi = \sqrt{1 + |\nabla \eta|^2} \cdot \partial_n \Psi \Big|_{y = \eta(t, x)}$$

Note that

$$\Sigma_t = \{(x, y) : y - \eta(x) = 0\}.$$

So outward normal is $(1, -\nabla\eta)$.

Hence

$$\begin{aligned} [G(\eta)\psi](t, x) &= \nabla_{x, y} \psi \cdot (1, -\nabla\eta) \\ &= \partial_y \psi - \nabla_x \psi \cdot \nabla\eta \\ &= \partial_t \eta. \end{aligned}$$

Also,

$$\left. \begin{aligned} \partial_y \psi &= [G(\eta)\psi] + \nabla\psi \cdot \nabla\eta \\ \nabla_x \psi &= \nabla\psi - (\partial_y \psi) \nabla\eta \end{aligned} \right\} \Rightarrow \partial_y \psi = \frac{[G(\eta)\psi] + \nabla\psi \cdot \nabla\psi}{1 + |\nabla\eta|^2} := \mathcal{B}$$

Thus we have

$$\partial_y \psi := \mathcal{B}$$

$$\nabla_x \psi := \nabla\psi - \mathcal{B} \nabla\eta := \mathcal{V}$$

Remark. Later the unknown $\psi - T_{\mathcal{B}} \eta$ will turn out to be important.

Further, we see now that (η, ψ) solve the Cauchy problem

$$\begin{cases} \eta(0, x) = \eta_0, \quad \psi(0, x) = \psi_0 \\ \partial_t \eta - (G(\eta)\psi) = 0 \\ \partial_t \psi + g\eta - H(\eta) + \frac{1}{2} |\nabla \psi|^2 - \frac{1}{2} \frac{G(\eta)\psi + \nabla \eta \cdot \nabla \psi}{1 + |\nabla \eta|^2} = 0. \end{cases}$$

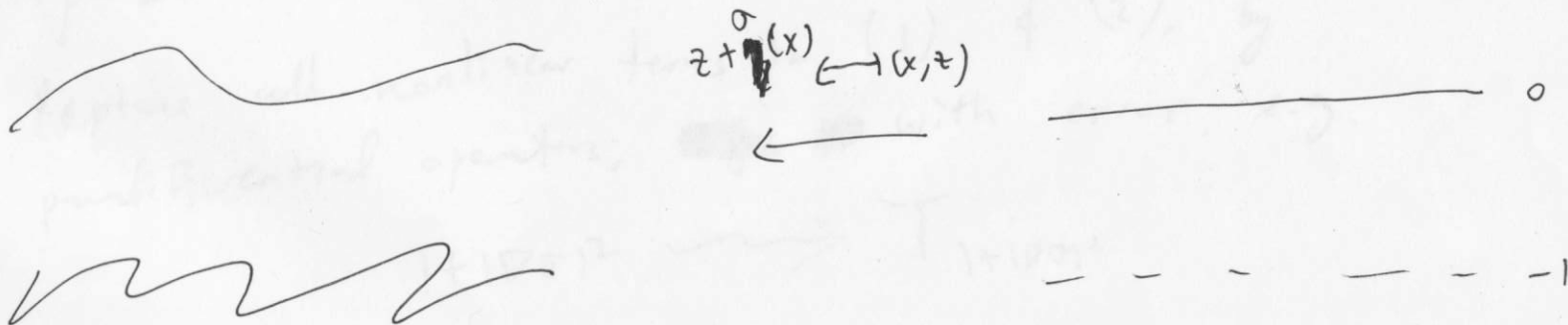
Idea: Paralinearize \rightsquigarrow reduce to simple hyperbolic form.
(paradifferential symmetrizer).

- Existence, uniqueness, regularity $(\eta_0, \psi_0) \in H^{s+1/2} \times H^s$
 $\Rightarrow (\eta, \psi) \in H^{s+1/2} \times H^s$ $d\eta, \psi > 2 + d/2$.
- Smoothing effect.

Full paralinearization: (will follow "Parad. of DN operator, and regularity of 3D WW" by Alazard, Métivier)

$$G(\sigma) \psi = \underbrace{Op(\Lambda_\sigma)}_{\text{explicit}} \psi + \underbrace{B(\sigma, \psi)}_{\text{more regular than } \sigma, \psi} + R(\sigma, \psi).$$

Change of variables.



Define

$$v(x, z) = \psi(x, z + \sigma(x))$$

$$v|_{z=0} = \psi|_{y=\sigma(x)} = \psi.$$

and

$$(1 + |\nabla\sigma|^2) \partial_z^2 v + \Delta v - 2 \nabla\sigma \cdot \nabla \partial_z v - \partial_z v \Delta\sigma = 0 \quad (1)$$

Since

$$\partial_z v = \partial_y \psi$$

$$\nabla_x v = \nabla_x \psi + (\partial_y \psi) \nabla\sigma$$

$$G(\sigma) \psi = \partial_y \psi - \nabla_x \psi \cdot \nabla\sigma$$

$$= \partial_z v - (\nabla_x v - (\partial_z v) \nabla\sigma) \cdot \nabla\sigma$$

$$G(\sigma) \psi = (1 + |\nabla\sigma|^2) \partial_z v - \nabla\sigma \cdot \nabla v |_{z=0} \quad (2)$$

Idea for penalization:

Replace all nonlinear terms in (1) & (2), by pseudo-differential operators, ~~with~~ with error, e.g.

$$1 + |\nabla\sigma|^2 \rightsquigarrow T_{1 + |\nabla\sigma|^2}$$